## Chapter (1)

## Force system

## 1. Analytical method

The magnitude and direction of the resultant force may be obtained, analytically, as discussed below:

(a)

$F_{3}$
Free body diagram
(b)

Fig. 1

1. First of ahl analysis all forces into two components in free body diagram as shown in Fig. 1 (b)
2. Resolve the forces horizontally and vertically and find their sums, i.e. $\sum F_{x}$ and $\sum F_{y}$. We know that

Sum of horizontal components of the forces at x -axis,
$\sum F_{x}=F_{1} \cos \theta_{1}-F_{2} \cos \theta_{2}$
Sum of vertically components of the forces at y -axis,
$\sum F_{y}=F_{1} \sin \theta_{1}+F_{2} \sin \theta_{2}-F_{3}$
3. Magnitude of the resultant force,
$R=\sqrt{\left(\sum F_{x}\right)^{2}+\left(\sum F_{y}\right)^{2}}$
4. If $\theta_{R}$ is the angle, which the resultant force makes with the horizontal, then
$\theta_{R}=\tan ^{-1}\left(\sum F_{y} / \sum F_{x}\right)$

## 2. Graphical method

The magnitude and position of the resultant force may also be obtained graphically as discussed below:

1. First of all, draw the space diagram with the positions of the several Forces, as shown in Fig. 1 (a).
2. Select suitable scale for all forces.
3. Draw the force diagram as shown in Fig. 2


Fig. 2: Force diagram
3. Determine the resultant force by using equation (5)
$R=\overline{o c} \times$ scale
4. The direction of resultant force can be estimated from Fig.2.

## Example: 1

The screw eye in Fig. 3 is subjected to two forces, $\boldsymbol{F}_{1}$ and $\boldsymbol{F}_{2}$.
Determine the magnitude and direction of the resultant force.


Fig. 3: Screw eye.

## Solution:

## 1. Analytical method

Draw free body diagram and analysis all forces into two components in x and y directions.


Free body diagram
From free body diagram we have;
$\sum F_{x}=F_{1} \cos \theta_{1}+F_{2} \cos \theta_{2}=100 \cos 15^{\circ}+150 \cos 80^{\circ}=122.63$
$\sum F_{y}=F_{1} \sin \theta_{1}+F_{2} \sin \theta_{2}=100 \sin 15^{\circ}+150 \sin 80^{\circ}=173.60 N$
The resultant force can be calculated by;
$R=\sqrt{\left(\sum F_{x}\right)^{2}+\left(\sum F_{y}\right)^{2}}=\sqrt{(122.63)^{2}+(173.60)^{2}}=213 \mathrm{~N}$
The direction of resultant force can be calculated by;
$\theta_{R}=\tan ^{-1} \frac{\sum F_{y}}{\sum F_{x}}=\tan ^{-1} \frac{173.60}{122.63}=54.8^{\circ}$

## 2. Graphical method

1. First of all, draw the space diagram with the positions of the several Forces, as shown in Fig. 3.
2. Select suitable scale for all forces.

Take allowable scale for all $1 \mathrm{~cm}=25 \mathrm{~N}$
3. Draw the force diagram as shown in Fig.4.


Fig. 4
From Fig. 4;

$$
R=o b \times \text { scale }=8.4 \times 25=210 \mathrm{~N}
$$

And

$$
\theta_{R}=55^{\circ}
$$

## Example: 2

The force $\mathrm{F}=450 \mathrm{~N}$ acts on the frame shown in Fig. 5. Resolve this force into components acting along members AB and AC , and determine the magnitude of each component.


Fig. 5:

## Solution:

## 1. Analytical method

Draw free body diagram and analysis all forces into two components in x and y directions.


Free body diagram

By using the Equilibrium conditions to obtain the forces $\boldsymbol{F}_{\boldsymbol{A} C}$ and $\boldsymbol{F}_{\boldsymbol{A} \boldsymbol{B}}$.

1. $\sum F_{x}=0$
$\therefore F_{A B} \cos 45-F_{A C} \cos 30=0$
$\therefore F_{A B}=F_{A C} \frac{\cos 30}{\cos 45}$
2. $\sum F_{y}=0$
$\therefore F_{A B} \sin 45-F_{A C} \sin 30-450=0$
$\therefore F_{A B} \sin 45-F_{A C} \sin 30=450$
By solving Eqns. (1) and (2) we have
$F_{A C} \frac{\cos 30}{\cos 45} \sin 45-F_{A C} \sin 30=450$
$F_{A C} \frac{\sin 45}{\cos 45} \cos 30-F_{A C} \sin 30=450$
$F_{A C} \tan 45 \cos 30-F_{A C} \sin 30=450$
$F_{A C}[\tan 45 \cos 30-\sin 30]=450$
$\therefore F_{A C}=\frac{450}{[\tan 45 \cos 30-\sin 30]}=1229.4 N$ Tension force
And
$F_{A B}=F_{A C} \frac{\cos 30}{\cos 45}=1229.4 \frac{\cos 30}{\cos 45}=1505.7 \mathrm{~N}$ Compression force

## Example: 3

Determine the magnitude of the resultant force and its direction measured counterclockwise from the positive $\boldsymbol{x}$ axis. Using analytical and graphical methods.


## Solution:

## I. Analytical Method

From the free body diagram, we get;
$\sum F_{x}=500 \cos 30^{\circ}-700 \cos 15^{\circ}=-243.14 N$
$\sum F_{y}=500 \sin 30^{\circ}-700 \sin 15^{\circ}=+68.83 N$
The resultant force can be given as
$\boldsymbol{R}=\sqrt{\left(\sum \boldsymbol{F}_{x}\right)^{2}+\left(\sum \boldsymbol{F}_{y}\right)^{2}}=$
$\sqrt{(-243.14)^{2}+(+68.83)^{2}}=252.7 \mathrm{~N}$
The angle of the resultant force can be given using the Equation;


Free body diagram
$\varphi=\tan ^{-1} \frac{\sum F_{y}}{\sum \bar{F}_{x}}=\tan ^{-1} \frac{+68.83}{-243.14}=15.8^{\circ}$
$\therefore \theta=180-\varphi=180^{\circ}-15.8^{o}=164.19^{\circ}$

## II. Graphical Method

From the force polygon diagram, we get out;
The resultant force can be measured from the diagram
$R=260 N$ and $\theta=165^{\circ}$


## Example: 4

Determine the resultant $\mathbf{R}$ of the three tension forces acting on the eye bolt. Find the magnitude of $R$ and the angle which $R$ makes with the positive $x$-axis.


## Answer:

## a. Analytical Method

From the free body diagram, we have


Free body diagram
$\sum F_{x}=20 \sin 30^{\circ}+8 \cos 45^{\circ}=15.656 k N$
$\sum F_{y}=20 \cos 30^{\circ}-4-8 \sin 45^{\circ}=7.664 k N$
$R=\sqrt{\sum{F_{x}}^{2}+\sum{F_{y}}^{2}}=\sqrt{(15.656)^{2}+(7.664)^{2}}=17.43 \mathrm{kN}$
$\theta=\tan ^{-1} \frac{\sum F_{y}}{\sum F_{x}}=\tan ^{-1} \frac{7.664}{15.656}=26.08^{\circ}$

## b. Graphical Method

From force diagram we get;
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Force diagram
$\therefore R=\overline{o c} \times$ scale $=8.7 \times 2=17.4 \mathrm{kN}$
$\therefore \theta_{R}=26^{\circ}$

## Example: 5

Determine the resultant $\mathbf{R}$ of the two forces applied to the bracket. Write $\mathbf{R}$ in terms of unit vectors along the $x$ - and $y$-axes shown.


## Solution:

## I. Analytical Method

From the free body diagram, we get;
$\sum F_{x}=200 \cos 35^{\circ}-150 \cos 60^{\circ}$

$$
=88.8304 \mathrm{~N}
$$

$\sum F_{y}=200 \sin 35^{\circ}+150 \sin 60^{\circ}$
$244.62 N$
The resultant force can be given as
$R=\sqrt{\left(\sum F_{x}\right)^{2}+\left(\Sigma F_{y}\right)^{2}}=$
$\sqrt{(88.8304)^{2}+(244.62)^{2}}=260.25 N$
The angle of the resultant force can be given using the Equation;
$\varphi=\tan ^{-1} \frac{\sum F_{y}}{\sum \bar{F}_{x}}=\tan ^{-1} \frac{244.62}{88.8304}=70.04^{\circ}$
$\therefore \theta=\boldsymbol{\varphi}=70.04{ }^{0}$

## II. Graphical Method

From force diagram we get;


Force Polygon diagram
$\therefore R=\overline{o b} \times$ scale $=10.3 \times 25=257.5 \mathrm{kN}$
$\therefore \theta_{R}=70^{\circ}$

## Example: 6

At what angle must the $400 \mathbf{N}$ force be applied in order that the resultant $\mathbf{R}$ of the two forces have a magnitude of $\mathbf{1 0 0 0} \mathbf{N}$ ? For this condition what will be the angle between $\mathbf{R}$ and the horizontal?


## Solution:

## Analytical Method

From the free body diagram, we get;
The resultant force can be given as


Free body diagram
$\therefore \cos \theta=\frac{(1000)^{2}-(400)^{2}-(700)^{2}}{(2 \times 400 \times 700)}$
$\therefore \boldsymbol{\theta}=\cos ^{-1}\left[\frac{(\mathbf{1 0 0 0})^{2}-(400)^{2}-(700)^{2}}{(2 \times 400 \times 700)}\right]=51.31^{\circ}$
$\sum F_{x}=-400 \cos \theta^{0}-700=-950.04 N$
$\sum F_{y}=400 \sin \theta^{o}=312.22 N$
The angle of the resultant force can be given using the Equation;
$\beta=\tan ^{-1} \frac{\sum F_{y}}{\sum F_{x}}=\tan ^{-1} \frac{312.22}{-950.04}=18.19^{\circ}$

## FUNDAMENTAL PROBLEMS

1. Determine the magnitude of the resultant force acting on the screw eye and its direction measured clockwise from the x axis.

2. Two forces act on the hook. Determine the magnitude of the resultant force.

3. The vertical force F acts downward at A on the two membered frame. Determine the magnitudes of the two components of F directed along the axes of $A B$ and $A C$ Set $F=500 \mathrm{~N}$.

4. Solve Prob. 3 with $\mathrm{F}=350 \mathrm{~N}$.

## Chapter (2)

## Moment of a force

## Objective

■ To discuss the concept of the moment of a force and show how to calculate it in two and three dimensions.

■ To provide a method for finding the moment of a force about a specified axis.

## Introduction

When a force is applied to a body it will produce a tendency for the body to rotate about a point that is not on the line of action of the force. This tendency to rotate is sometimes called a torque, but most often it is catled the moment of a force or simply the moment. For example, consider a wrench used to unscrew the bolt in Fig. 6 (a). If a force is applied to the handle of the wrench it will tend to turn the bolt about point $O$ (or the zaxis). The magnitude of the moment is directly proportional to the magnitude of $\mathbf{F}$ and the perpendicular distance or moment arm $d$. The larger the force or the longer the moment arm, the greater the moment or turning effect. Note that if the force $\mathbf{F}$ is applied at an angle $\theta \neq$ $90^{\circ}$, Fig. $6(b)$, then it will be more difficult to turn the bolt since the moment arm $d^{\prime}=d \sin \theta$ will be smaller than $d$. If $\mathbf{F}$ is applied along the wrench, Fig. 6 (c), its moment arm will be zero since the line of action of $\mathbf{F}$ will intersect point $O$ (the $z$ axis). As a result, the moment of $\mathbf{F}$ about $O$ is also zero and no turning can occur.


Fig. 6:
We can generalize the above discussion and consider the force $\mathbf{F}$ and point $O$ which lie in the shaded plane as shown in Fig. 7 (a). The moment Mo about
point $O$, or about an axis passing through $O$ and perpendicular to the plane, is a vector quantity since it has a specified magnitude and direction.

(b)
$M_{O}=F \times d$
Where $\boldsymbol{d}$ is the moment arm or perpendicular distance from the axis at point $\boldsymbol{O}$ to the line of action of the force. Units of moment magnitude consist of force times distance, e.g., $N . m$ or $\mathbf{l b}$, $f t$.

Direction. The direction of $\boldsymbol{M}_{\boldsymbol{o}}$ is defined by its moment axis, which is perpendicula to the plane that contains the force $\mathbf{F}$ and its moment arm $\boldsymbol{d}$. The right-hand rule is used to establish the sense of direction of Mo. According to this rule, the natural curl of the fingers of the right hand, as they are drawn towards the palm, represent the rotation, or if no movement is possible, there is a tendency for rotation caused by the moment. As this action is performed, the thumb of the right hand will give the directional sense of Mo, Fig. 7 (a). Notice that the moment vector is represented three-dimensionally by a curl around an arrow. In two dimensions this vector is represented only by the curl as in Fig. 7 (b). Since in this case the moment will tend to cause a counterclockwise rotation, the moment vector is actually directed out of the page.

Resultant Moment. For two-dimensional problems, where all the forces lie within the $\boldsymbol{x}-\boldsymbol{y}$ plane, Fig. 8, the resultant moment $(\mathbf{M} R) O$ about point $\boldsymbol{O}$ (the $\boldsymbol{z}$ axis) can be determined by finding the algebraic sum of the moments caused
by all the forces in the system. As a convention, we will generally consider positive moments as counterclockwise since they are directed along the positive $z$ axis (out of the page). Clockwise moments will be negative. Doing this, the directional sense of each moment can be represented by a plus or minus sign. Using this sign convention, the resultant moment in Fig. 8 is therefo
$\left(+\left(M_{R}\right)_{o}=\sum F d ;\right.$

$$
\left(M_{R}\right)_{o}=F_{1} d_{1}-F_{2} d_{2}+F_{3} d_{3}
$$



Fig. 8:
If the numerical result of this sum is a positive scalar, $\left(\boldsymbol{M}_{\boldsymbol{R}}\right)_{o}$ will be a counterclockwise moment (out of the page); and if the result is negative, $\left(\boldsymbol{M}_{\boldsymbol{R}}\right)_{o}$ will be a clockwise moment (into the page).

## Example: 1

For each case illustratedin Fig. 9, determine the moment of the force about


Fig. 9

## Solution:

The line of action of each force is extended as a dashed line in order to establish the moment arm $\boldsymbol{d}$. Also illustrated is the tendency of rotation of the member as caused by the force. Furthermore, the orbit of the force about $\boldsymbol{O}$ is shown as a colored curl. Thus,

Fig. 9 (a) $\quad M_{O}=(100 N)(2 m)=200$ N.m $)$
Fig. 9 (b)

$$
\left.M_{O}=(50 \mathrm{~N})(0.75 \mathrm{~m})=37.5 \mathrm{~N} . \mathrm{m}\right)
$$

Fig. 9 (c)

$$
\left.M_{O}=(40 \mathrm{Ib})\left(4 f t+2 \cos 30^{\circ} f t\right)=229 \mathrm{Ib} . f t\right)
$$

Fig. 9 (d)

$$
\left.M_{O}=(60 \mathrm{Ib})\left(1 \sin 45^{\circ} \mathrm{ft}\right)=42.4 \mathrm{Ib} \cdot f t^{\}}\right)
$$

Fig. 9 (b)

$$
M_{O}=(7 k N)(4 m-1 m)=21.0 k N \cdot m
$$

## Example: 2

Determine the resultant moment of the four forces acting on the rod shown in Fig. 10 about point $\boldsymbol{O}$.


Fig. 10

## Solution:

Assuming that positive moments act in the $+\mathbf{k}$ direction, i.e., counterclockwise, we have

$$
\begin{aligned}
& \left(+\left(M_{R}\right)_{o}=\sum F d\right. \\
& \quad\left(M_{R}\right)_{o}=-50 N(2 m)+60 N(0)+20 N\left(3 \sin 30^{\circ} m\right)-40 N(4 m+ \\
& \left.3 \cos 30^{\circ} m\right)
\end{aligned}
$$

$$
\left.\therefore\left(M_{R}\right)_{o}=-334 \mathrm{~N} . \mathrm{m}=334 \mathrm{~N} . \mathrm{m}\right)
$$

For this calculation, note how the moment-arm distances for the $20-\mathrm{N}$ and $40-$ N forces are established from the extended (dashed) lines of action of each of these forces.

## Example: 3

Calculate the magnitude of the moment about the base point $O$ of the $600-\mathrm{N}$ force in five different ways.

## Solution:

(I) The moment arm to the $600-\mathrm{N}$ force is

$$
\mathrm{d}=4 \cos 40^{\circ}+2 \sin 40^{\circ}=4.35 \mathrm{~m}
$$

(1) By $M=F d$ the moment is clockwise and has the magnitude

$$
M_{O}=600(4.35)=2610 \mathrm{~N} \cdot \mathrm{~m} \text { Ans. }
$$


(II) Replace the force by its rectangular components at $A$, $F_{1}=600 \cos 40^{\circ}=460 \mathrm{~N}, \quad F_{2}=600 \sin 40^{\circ}=386 \mathrm{~N}$ By Avignôn's theorem, the moment becomes
(2)

$$
M_{O}=460(4)+386(2)=2610 \mathrm{~N} . \mathrm{m}
$$

Ans.


## Example: 4

Calculate the moment of the $250-\mathrm{N}$ force on the handle of the monkey wrench about the center of the bolt.


## Solution:


$250 \cos 15^{\circ}$
$\left.M=\left(250 \sin 15^{\circ} \times 30\right)-\left(250 \cos 15^{\circ} \times 200\right)=-46867.66 \mathrm{~N} . \mathrm{m}^{3}\right)$
$\times M=46867.66 \mathrm{~N} . \mathrm{m})$

## FUNDAMENTAL PROBLEMS

1. Determine the moment of the force about point $O$.


## Chapter (3)

## Equilibrium of a Rigid Rody

## 1. OBJECTIVES

■ To develop the equations of equilibrium for a rigid body.
■ To introduce the concept of the free-body diagram for a rigid body.
■ To show how to solve rigid-body equilibrium problems using the equations of equilibrium.

## 2. Conditions for Rigid-Body Equilibrium

In this section, we will develop both the necessary and sufficient conditions for the equilibrium of the rigid body in Fig. 11. As shown, this body is subjected to an external force and couple moment system that is the result of the effects of gravitational, electrical, magnetic, or contact forces caused by adjacent bodies. The internal forces caused by interactions between particles within the body are not shown in this figure because these forces occur in equal but opposite collinear pairs and hence will cancel out, a consequence of Newton's third law.


Fig. 11

## 3. Free-Body Diagrams

Successful application of the equations of equilibrium requires a complete specification of all the known and unknown external forces that act on the body. The best way to account for these forces is to draw a free-body diagram. This diagram is a sketch of the outlined shape of the body, which represents it as being isolated or "free" from its surroundings, i.e., a "free body." On this sketch it is necessary to show all the forces and couple moments that the surroundings exert on the body so that these effects can be accounted for when the equations
of equilibrium are applied. A thorough understanding of how to draw a freebody diagram is of primary importance for solving problems in mechanics.

## 4. Support Reactions

Before presenting a formal procedure as to how to draw a free-body diagram, we will first consider the various types of reactions that occur at supports and points of contact between bodies subjected to coplanar force systems. As a general rule,

- If a support prevents the translation of a body in a given direction, then a force is developed on the body in that direction.
- If rotation is prevented, a couple moment is exerted on the body

For example, let us consider three ways in which a horizontal member, such as a beam, is supported at its end. One method consists of a rolter or cylinder, Fig. 12 (a). Since this support only prevents the beam from translating in the vertical direction, the roller will only exert a force on the beam in this direction, Fig. 12 (b).

The beam can be supported in a more restrictive manner by using a pin, Fig. 12 (c). The pin passes through a hole in the beam and two leaves which are fixed to the ground. Here the pin can prevent translation of the beam in any direction f, Fig. $12(d)$, and so the pin must exert a force $\mathbf{F}$ on the beam in this direction. For purposes of analysis, it is generally easier to represent this resultant force $\mathbf{F}$ by its two rectangular components $\mathbf{F}_{x}$ and $\mathbf{F}_{y}$, Fig. 12 (e). If $\boldsymbol{F}_{x}$ and $\boldsymbol{F}_{y}$ are known, then $\boldsymbol{F}$ and f can be calculated.

The most restrictive way to support the beam would be to use a fixed support as shown in Fig. 12 (\%), This support will prevent both translation and rotation of the beam. To do this a force and couple moment must be developed on the beam at its point of connection, Fig. $12(\mathrm{~g})$. As in the case of the pin, the force is usually represented by its rectangular components $\boldsymbol{F}_{\boldsymbol{x}}$ and $\boldsymbol{F}_{\boldsymbol{y}}$.

Table Nists other common types of supports for bodies subjected to coplanar force systems. (In all cases the angle $u$ is assumed to be known.) Carefully study each of the symbols used to represent these supports and the types of reactions they exert on their contacting members.


Fig. 12.

| Types of Connection Reaction | Number of Unknowns |
| :---: | :---: |
| (1) <br> cable | One unknown. The reaction is a tension force which acts away from the member in the direction of the cable. |
| (2) <br> or <br> weightless link | One unknown. The reaction is a force which acts along the axis of the link. |
| (3) | One unknown. The reaction is a force which acts perpendicular to the surface at the point of contact. |
| (4) | One unknown. The reaction is a force which acts perpendicular to the surface at the point of contact. |
|  | One unknown. The reaction is a force which acts perpendicular to the surface at the point of contact. |
|  | One unknown. The reaction is a force which acts perpendicular to the slot. |
|  | One unknown. The reaction is a force which acts perpendicular to the rod. |

## 5. Equations of Equilibrium

In Sec. 2 we developed the two equations which are both necessary and sufficient for the equilibrium of a rigid body, namely, $\sum \mathbf{F}=\mathbf{0}$ and $\sum \mathbf{M}_{o}=\mathbf{0}$. When the body is subjected to a system of forces, which all lie in the $x-y$ plane, then the forces can be resolved into their $\boldsymbol{x}$ and $\boldsymbol{y}$ components. Consequently, the conditions for equilibrium in two dimensions are
$\sum F_{x}=0$
$\sum F_{y}=0$
$\sum M_{O}=0$

## Example: 1

Draw the free-body diagram of the uniform beam shown in Fig. 13. The beam has a mass of 100 kg . And determine the reactions moment and forces.


Fig. 13.

## Solution:

The free-body diagram of the beam is shown in Fig. 14. Since the support at $\boldsymbol{A}$ is fixed, the wall exerts three reactions on the beam, denoted as $\boldsymbol{A}_{\boldsymbol{x}}, \boldsymbol{A}_{\boldsymbol{y}}$, and $\boldsymbol{M}_{\boldsymbol{A}}$. The magnitudes of these reactions are anknown, and their sense has been assumed. The weight of the beam, $\mathrm{W}=100(9.81) \mathrm{N}=\mathbf{9 8 1} \mathrm{N}$, acts through the beam's center of gravity $\boldsymbol{G}$, which is $\mathbf{3} \boldsymbol{m}$ from $\boldsymbol{A}$ since the beam is uniform.


Effect of gravity (weight) acting on beam

Fig. 14.
From the equilibrium conditions we get;

$$
\begin{array}{ll}
\sum F_{x}=0 & \therefore A_{x}=0 \\
\sum F_{y}=0 & \therefore A_{y}=1200+981=2181 N \\
\sum M_{O}=0 & \therefore M_{A}=(981 N \times 3 \mathrm{~m})+(1200 N \times 2 \mathrm{~m})=5343 \mathrm{~N} . \mathrm{m}
\end{array}
$$

## Example: 2

Determine the magnitude $\boldsymbol{T}$ of the tension in the supporting cable and the magnitude of the force on the pin at $\boldsymbol{A}$ for the jib crane shown. The beam $\boldsymbol{A B}$ is a standard $\mathbf{0 . 5 - m}$ I-beam with a mass of 95 kg per meter of length.

## Solution:

## Algebraic solution.

The system is symmetrical about the vertical $x-y$ plane through the center of the beam, so
 the problem may be analyzed as the equilibrium of a coplanar force system. The free-body diagram of the beam is shown in the figure with the pin reaction at $\boldsymbol{A}$ represented in terms of its two rectangular components. The weight of the beam is $\mathbf{9 5 ( 1 0} \mathbf{~} \mathbf{~}^{(5) 9.81=4.66} \mathbf{k N}$ and acts through its center. Note that there are three unknowns $\boldsymbol{A}_{\boldsymbol{x}}, \boldsymbol{A}_{\boldsymbol{y}}$, and $\boldsymbol{T}$, which may be found from the three equations of equilibrium. We begin with a moment equation about $\boldsymbol{A}$, which eliminates two of the three unknowns from the equation. In applying the moment equation about $A$, it is simpler to consider the moments of the $\boldsymbol{x}$ - and $\boldsymbol{y}$-components of $\mathbf{T}$ than it is to compute the perpendicular distance from $\mathbf{T}$ to $\boldsymbol{A}$. Hence, with the counterclockwise sense as positive we write
$\left[\sum M_{A}=0\right] \quad\left(T \cos 25^{\circ}\right) 0.25+\left(T \sin 25^{\circ}\right)(5-0.12)-10(5-1.5-$ $0.12)-4.66(2.5-0.12)=0$

From which $\quad T=19.61 \mathrm{kN}$
Equating the sums of forces in the $x$ - and $y$-directions to zero gives

| $\left[\Sigma F_{x}=0\right]$ | $A_{x}-19.61 \cos 25^{\circ}=0 \quad \therefore A_{x}=17.77 \mathrm{kN}$ |
| :--- | :--- |
| $\left[\sum F_{y}=0\right]$ | $A_{y}+19.61 \sin 25^{\circ}-466-10=0 \quad \therefore A_{y}=6.37 \mathrm{kN}$ |
| $\left\|A=\sqrt{A_{x}{ }^{2}+{A_{y}}^{2}}\right\|$ | $A=\sqrt{(17.77)^{2}+(6.37)^{2}}=18.88 \mathrm{kN} \quad$ Ans. |



Free-body diagram

## Graphical solution.

The principle that three forces in equilibrium must be concurrent is utilized for a graphical solution by combining the two known vertical forces of 4.66 and 10 kN into a single $14.66-\mathrm{kN}$ force, located as shown on the modified free-body diagram of the beam in the lower figure. The position of this resultant load may easily be determined graphically or algebraically. The intersection of the 14.66$\boldsymbol{k N}$ force with the line of action of the unknown tension $\mathbf{T}$ defines the point of concurrency $\boldsymbol{O}$ through which the pin reaction $\boldsymbol{A}$ pust pass. The unknown magnitudes of $\mathbf{T}$ and $\mathbf{A}$ may now be found by adding the forces head-to-tail to form the closed equilibrium polygon of forces, thus satisfying their zero-vector sum. After the known vertical load is land off to a convenient scale, as shown in the lower part of the figure, a line representing the given direction of the tension T is drawn through the the of the $14.66-k N$ vector. Likewise, a line representing the direction of the pin reaction $\mathbf{A}$, determined from the concurrency established with the free-body diagram, is drawn through the tail of the $\mathbf{1 4 . 6 6 - k N}$ vector. The intersection of the lines representing vectors $\mathbf{T}$ and A establishes the magnitudes $T$ and $A$ necessary to make the vector sum of the forces equal to zero. These magnitudes are scaled from the diagram. The $\mathbf{x}$ - and $\boldsymbol{y}$-components of $\mathbf{A}$ may be constructed on the force polygon if desired.


## Example: 3

The uniform beam has a mass of $50 \mathbf{k g}$ per meter of length. Determine the reactions at the supports.


## Solution:

From the free-body diagram we get

$$
\begin{array}{ll}
{\left[\sum F_{x}=0\right] \quad \therefore R_{B_{x}}=0} \\
{\left[\sum F_{y}=0\right] \quad \therefore R_{A_{y}}+R_{B_{y}}=300+180=480 \mathrm{kN}} \\
{\left[\sum M_{A}=0\right]} & (300 \times 2.4)+(180 \times 1.8)-\left(R_{B_{y}} \times 3.6\right)=0 \\
& R_{B_{y}}=\frac{(300 \times 2.4)+(180 \times 1.8)}{3.6}=290 \mathrm{kN}
\end{array}
$$

From Eqn. (1) $\quad \therefore R_{A_{y}}=480-290=190 k N$


Free-body diagram

## Example: 4

The $500-\mathrm{kg}$ uniform beam is subjected to the three external loads shown. Compute the reactions at the support point $O$. The $x-y$ plane is vertical.


## Solution:

From the free-body diagram we get


Free-body diagram
$\left[\Sigma F_{x}=0\right] \quad R_{O_{x}}-1.36=0 \quad \therefore R_{O_{x}}=1.36 \mathrm{kN}$
$\left[\Sigma F_{y}=0\right] \quad \therefore R_{0_{y}}+1.4-5-2.67=0 \quad \therefore R_{O_{y}}=6.27 \mathrm{kN}$
$\left[\Sigma M_{0}=0\right] \quad(2.67 \times 4.8)+(5 \times 2.4)-15-(1.4 \times 1.2)-M_{0}=0$
$\therefore M_{0}=8.136 \mathrm{kN} . \mathrm{m}$ CCW

## Example: 5

Determine the horizontal and vertical components of reaction on the beam caused by the pin at $B$ and the rocker at $A$ as shown in Fig.. Neglect the weight of the beam.


## Solution:

From the free-body diagram we get


Equations of Equilibrium. Summing forces in the $x$ direction yields
$\xrightarrow{+} \sum F_{x}=0 ; \quad 600 \cos 45^{\circ}-R_{B_{x}}=0$

$$
\therefore R_{B_{x}}=424 \mathrm{~N}
$$

Ans.
A direct solution for $\boldsymbol{A}_{\boldsymbol{y}}$ can be obtained by applying the moment equation $+\cup \sum M_{B}=0$ about point $\boldsymbol{B}$.
$v+\sum M_{B}=0 ;$

$$
\begin{gathered}
(100 \times 2)+\left(600 \sin 45^{\circ} \times 5\right)-\left(600 \cos 45^{\circ} \times 2\right)-\left(R_{A_{y}} \times 7\right)=0 \\
\therefore R_{A_{y}}=319 N
\end{gathered} \text { Ans. }
$$

Summing forces in the $y$ direction, using this result, gives

$$
\begin{array}{r}
+\uparrow \sum F_{y}=0 ; \quad 319-600 \sin 45^{\circ}-100-200+R_{B_{y}}=0 \\
\therefore R_{B_{y}}=405 \mathrm{~N} \tag{Ans}
\end{array}
$$

## FUNDAMENTAL PROBLEMS

1. Draw the free-body diagram of the dumpster $\boldsymbol{D}$ of the truck, which has a mass of $2.5 \mathbf{M g}$ and a center of gravity at $\boldsymbol{G}$. It is supported by a pin at $\boldsymbol{A}$ and a pin-connected hydraulic cylinder $\boldsymbol{B C}$ (short link). Determine the horizontal and vertical components

2. Determine the support reactions on the member. The collar at $A$ is fixed to the member and can slide vertically along the vertical shaft.

3. Determine the horizontal and vertical components of reaction on the beam caused by the pin at $B$ and the rocker at $A$ as shown in Figure. Neglect the weight of the beam.

4. The jib crane is pin connected at $\boldsymbol{A}$ and supported by a smooth collar at $\boldsymbol{B}$. Determine the roller placement $\boldsymbol{x}$ of the $\mathbf{5 0 0 0} \mathbf{l b}$ load so that it gives the
maximum and minimum reactions at the supports. Calculate these reactions in each case. Neglect the weight of the crane. Require $\mathbf{4 f t} \leq \boldsymbol{x} \leq \mathbf{1 0} \mathbf{f t}$.

5. The crane consists of three parts, which have weights of $W_{1}=3500 \mathrm{lb}, W 2$ $=900 \mathrm{lb}, W_{3}=1500 \mathrm{lb}$ and centers of gravity at $G_{1}, G_{2}$, and $G_{3}$, respectively. Neglecting the weight of the boom, determine (a) the reactions on each of the four tires if the load is hoisted at constant velocity and has a weight of 800 lb , and (b), with the boom held in the position shown, the maximum load the crane can lift without tipping over.

6. The cantilevered jib crane is used to support the load of 780 lb . If the trolley $\boldsymbol{T}$ can be placed anywhere between $1.5 \mathrm{ft} \leq x \leq 7.5$ $f t$, determine the maximum magnitude of reaction at the supports $\boldsymbol{A}$ and $\boldsymbol{B}$. Note that the supports are collars that allow the crane to rotate freely about the vertical axis. The collar at $\boldsymbol{B}$
 supports a force in the vertical direction, whereas the one at $\boldsymbol{A}$ does not.

## Chapter (4)

## Friction

## CHAPTER DRJECTIVES

- To introduce the concept of dry friction and show how to analyze the equilibrium of rigid bodies subjected to this force.
- To present specific applications of frictional force analysis on wedges, screws, belts, and bearings.
$\square$ To investigate the concept of rolling resistance


## 1. Characteristics of Dry Friction

Friction is a force that resists the movement of two contacting surfaces that slide relative to one another. This force always acts tangent to the surface at the points of contact and is directed so as to oppose the possible or existing motion between the surfaces.

In this chapter, we will study the effects of dry friction, which is sometimes called Coulomb friction since its characteristics were studied extensively by C. A. Coulomb in 1781. Dry friction occurs between the contacting surfaces of bodies when there is no lubricating fluid. *

## 2. Theory of Drantrition.

The theory of dry friction can be explained by considering the effects caused by pulling horizontally on a block of uniform weight $\mathbf{W}$ which is resting on a rough horizontal surface that is nonrigid or deformable, Fig. 1 (a). The upper portion of the block, however, can be considered rigid. As shown on the freebody diagram of the block, Fig. 1 (b), the floor exerts an uneven distribution of both normal force $\mathbf{\Delta} \boldsymbol{N}_{n}$ and frictional force $\mathbf{\Delta} \mathbf{F}_{n}$ along the contacting surface. For equiribrium, the normal forces must act upward to balance the block's weight $\mathbf{W}$, and the frictional forces act to the left to prevent the applied force $\mathbf{P}$ from moving the block to the right. Close examination of the contacting surfaces between the floor and block reveals how these frictional and normal forces develop, Fig. 1 (c). It can be seen that many microscopic irregularities exist between the two surfaces and, as a result, reactive forces $\Delta \mathbf{R}_{n}$ are developed at each point of contact. As shown, each reactive force contributes both a frictional component $\Delta \mathbf{F}_{n}$ and a normal component $\Delta \mathbf{N}_{n}$.


Fig. 1

Equilibrium. The effect of the distributed normal and frictional loadings is indicated by their resultants $\mathbf{N}$ and $\mathbf{F}$ on the free-body diagram, Fig. $1(d)$. Notice that $\mathbf{N}$ acts a distance $x$ to the right of the line of action of W, Fig. $1(d)$. This location, which coincides with the centroid or geometric center of the normal force distribution in Fig. 1(b), is necessary in order to balance the "tipping effect" caused by P. For example, if $\mathbf{P}$ is applied at a height $h$ from the surface, Fig. $1(d)$, then moment equilibrium about point $O$ is satisfied if $\boldsymbol{W}_{\boldsymbol{x}}=\boldsymbol{P h}$ or $\boldsymbol{x}=\boldsymbol{P} \boldsymbol{h} / \boldsymbol{W}$.

Impending Motion. In cases where the surfaces of contact are rather "slippery," the frictional force $\mathbf{F}$ may not be great enough to balance $\mathbf{P}$, and consequently the block will tend to slip. In other words, as $\boldsymbol{P}$ is slowly increased, $F$ correspondingly increases until it attains a certain maximum value $\boldsymbol{F}_{s}$, called the limiting static frictional force, Fig. 1 (e). When this value is reached, the block is in unstable equilibrium since any further increase in $\boldsymbol{P}$ will cause the block to move. Experimentally, it has been determined that this limiting static frictional force $\boldsymbol{F}_{s}$ is directly proportional to the resultant normal force N. Expressed mathematically,

(e)

Fig. 1 (cont.)
$F_{S}=\mu_{s} N$

Where the constant of proportionality, $\boldsymbol{\mu}_{s}(\mathrm{mu}$ "sub" $s$ ), is called the coefficient of static friction.
Thus, when the block is on the verge of sliding, the normal force $\mathbf{N}$ and frictional force $\boldsymbol{F}_{s}$ combine to create a resultant $\boldsymbol{R}_{s}$, Fig. $1(e)$. The angle $\boldsymbol{\phi}_{s}$ (phi "sub" $s$ ) that $\boldsymbol{R}_{s}$ makes with $\boldsymbol{N}$ is called the angle of static friction.
From the figure,
$\phi_{s}=\tan ^{-1}\left(\frac{F_{s}}{N}\right)=\tan ^{-1}\left(\frac{\mu_{s} N}{N}\right)=\tan ^{-1} \mu_{s}$

Typical values for $\boldsymbol{\mu}_{s}$ are given in Table 1. Note that these values can vary since experimental testing was done under variable conditions of coughness and cleanliness of the contacting surfaces. For applications, therefore, it is important that both caution and judgment be exercised when selecting a coefficient of friction for a given set of conditions. When a more accurate calculation of $\boldsymbol{F}_{\mathbf{s}}$ is required, the coefficient of friction should be determined directly by an experiment that involves the two materials to be used.

| Table 8-1 | Typical Values for $\mu_{s}$ |
| :--- | :---: |
| Contact <br> Materials | Coefficient of <br> Static Friction $\left(\mu_{\mathrm{s}}\right)$ |
| Metal on ice | $0.03-0.05$ |
| Wood on wood | $0.30-0.70$ |
| Leather on wood | $0.20-0.50$ |
| Leather on metal | $0.30-0.60$ |
| Aluminum on <br> aluminum | $1.10-1.70$ |

Motion. If the magnitude of $\boldsymbol{P}$ acting on the block is increased so that it becomes slightly greater than $\boldsymbol{F}_{s}$, the frictional force at the contacting surface will drop to a smaller value $\boldsymbol{F}_{\boldsymbol{k}}$, called the kinetic frictional force.
The block will begin to slide with increasing speed, Fig. 2 (a). As this occurs, the block will "ride" on top of these peaks at the points of contact, as shown in Fig. 2 (b). The continued breakdown of the surface is the dominant mechanism creating kinetic friction.

Experiments with sliding blocks indicate that the magnitude of the kinetic friction force is directly proportional to the magnitude of the resultant normal force, expressed mathematically as

(a)

(b)

Fig. 2
$F_{K}=\mu_{k} N$

Here the constant of proportionality, $\boldsymbol{\mu}_{\boldsymbol{k}}$, is called the coefficient of kinetic friction. Typical values for $m k$ are approximately 25 percent smaller than those listed in Table 1 for $\boldsymbol{\mu}_{s}$.
As shown in Fig. $2(a)$, in this case, the resultant force at the surface of contact, $\boldsymbol{R}_{\boldsymbol{k}}$, has a line of action defined by $\phi_{k}$. This angle is referred to as the angle of kinetic friction, where
$\phi_{k}=\tan ^{-1}\left(\frac{F_{k}}{N}\right)=\tan ^{-1}\left(\frac{\mu_{k} N}{N}\right)=\tan ^{-1} \mu_{k}$

By comparison, $\phi_{s} \geq \phi_{k}$.

The above effects regarding friction can be summarized by referring to the graph in Fig. 3, which shows the variation of the frictional force $\boldsymbol{F}$ versus the applied load $\boldsymbol{P}$. Here the frictional force is categorized in three different ways:


Fig. 3

- $\boldsymbol{F}$ is a static frictional force if equilibrium is maintained.
- $\boldsymbol{F}$ is a limiting static frictional force $\boldsymbol{F}_{s}$ when it reaches a maximum yalue needed to maintain equilibrium.
- $\boldsymbol{F}$ is a kinetic frictional force $\boldsymbol{F}_{\boldsymbol{k}}$ when sliding occurs at the contacting surface.

Notice also from the graph that for very large values of $\boldsymbol{P}$ or for high speeds, aerodynamic effects will cause $\boldsymbol{F}_{\boldsymbol{k}}$ and likewise $\boldsymbol{\mu}_{\boldsymbol{k}}$ to begin to decrease.

Characteristics of Dry Friction. As a result of experiments that pertain to the foregoing discussion, we can state the following rules which apply to bodies subjected to dry friction.

- The frictional force acts tangent to the contacting surfaces in a direction opposed to the motion or tendency for motion of one surface relative to another.
- The maximum static frictional foree $\boldsymbol{F}_{\boldsymbol{s}}$ that can be developed is independent of the area of contact, provided the normal pressure is not very low nor great enough to severely deform or crush the contacting surfaces of the bodies.
- The maximum static frictional force is generally greater than the kinetic frictional force for any two surfaces of contact. However, if one of the bodies is moving with a very low velocity oyer the surface of another, $\boldsymbol{F}_{\boldsymbol{k}}$ becomes approximately equal to $\boldsymbol{F}_{\boldsymbol{s}}$, i.e. $\boldsymbol{\mu}_{s} \approx \boldsymbol{\mu}_{k}$
- When stipping at the surface of contact is about to occur, the maximum static frictional force is proportional to the normal force, such that $\boldsymbol{F}_{s}=\boldsymbol{\mu}_{s} N$.
- When slipping at the surface of contact is occurring, the kinetic frictional force is proportional to the normal force, such that $F_{k}=\mu_{k} N$.


## Example: 1

Determine the maximum angle $\boldsymbol{\theta}$ which the adjustable incline may have with the horizontal before the block of mass $\boldsymbol{m}$ begins to slip. The coefficient of static friction between the block and the inclined surface is $\boldsymbol{\mu}_{s}$.


## Solution:

The free-body diagram of the block shows its weight $\boldsymbol{W}=\boldsymbol{m} \boldsymbol{g}$, the normal force $\boldsymbol{N}$, and the friction force $\boldsymbol{F}$ exerted by the incline on the block. The friction force acts in the direction to oppose the slipping which would occur if no friction were present.
Equilibrium in the $x$ - and $y$-directions requires

$\left[\sum F_{x}=0\right] \quad m g \sin \theta-F=0 \quad \therefore F=m g \sin \theta$
$\left[\Sigma F_{y}=0\right]-m g \cos \theta+N=0 \quad \therefore N=m g \cos \theta$
Dividing the first equation by the second gives $\boldsymbol{F} / \boldsymbol{N}=\boldsymbol{\operatorname { t a n }} \boldsymbol{\theta}$. Since the maximum angle occurs when $\boldsymbol{F}=\boldsymbol{F}_{\mathrm{max}}=\boldsymbol{\mu}_{s} \boldsymbol{N}$, for impending motion we have
$\mu_{s}=\tan \theta_{\max } \quad$ or $\quad \theta_{\max }=\tan ^{-1} \mu_{s}$

## Example: 2

Determine the range of values which the mass $\boldsymbol{m}_{0}$ may have so that the $\mathbf{1 0 0}$-kg block shown in the figure will neither start moving up the plane nor slip down the plane. The coefficient of static friction for the contact surfaces is $\mathbf{0 . 3 0}$.


## Solution:

The maximum value of $\boldsymbol{m}_{\boldsymbol{0}}$ will be given by the requirement for motion impending up the plane. The friction force on the block therefore acts down


Case I
the plane, as shown in the free-body diagram of the block for Case $I$ in the figure. With the weight $\boldsymbol{m g}=\mathbf{1 0 0}(9.81)=\mathbf{9 8 1} \mathrm{N}$, the equations of equilibrium give
$\left[\Sigma F_{y}=0\right] \quad N-m g \cos 20^{\circ}=0 \quad \therefore N=922 N$
$F_{\text {max }}=\mu_{s} N \quad \therefore F_{\text {max }}=0.30 \times 922=277 N$
$\left[\sum F_{x}=0\right] \quad m_{0} 9.81-277-981 \sin 20^{\circ}=0 \quad \therefore m_{0}=62.4 \mathrm{~kg} \quad$ Ans.

The minimum value of $\boldsymbol{m}_{0}$ is determined when motion is impending down the plane. The friction force on the block will act up the plane to oppose the tendency to move, as shown in the free-body diagram for Case II. Equilibrium in the $x$ direction requires


Case II

$$
\left[\sum F_{x}=0\right] \quad m_{0} 9.81+277-981 \sin 20^{\circ}=0 \quad \therefore m_{0}=6.01 \mathrm{~kg} \quad \text { Ans. }
$$

Thus, $\boldsymbol{m}_{0}$ may have any value from 6.01 to $62.4 \mathbf{k g}$, and the block will remain at rest.
In both cases equilibrium requires that the resultant of $\boldsymbol{F}_{\max }$ and $N$ be concurrent with the 981 N weight and the tension $T$.

## Example:

Determine the minimum force $\boldsymbol{P}$ to prevent the $30-k g$ rod $\boldsymbol{A B}$ from sliding. The contact surface at $\boldsymbol{B}$ is smooth, whereas the coefficient of static friction between the rod and the wall at $\boldsymbol{A}$ is $\boldsymbol{\mu}_{s}=\mathbf{0 . 2}$.

## Solution:



From free body diagram we get;

Bu using the equilibrium conditions

$$
\begin{aligned}
& 1-\sum M_{B}=0 \\
& \therefore\left(N_{A} \times 3\right)+(F \times 4)-(W \times 2)=0 \\
& F=\mu_{S} N_{A}=0.2 N_{A} \\
& \therefore\left(3 N_{A}\right)+\left(0.2 \times 4 N_{A}\right)=(2 W) \\
& \therefore N_{A}=\frac{2 W}{[3+(0.2 \times 4)]} \\
& \quad=\frac{2 \times 300 \times 9.81}{[3+(0.2 \times 4)]} \\
& \quad=154.89 \mathrm{~N}
\end{aligned} \quad \begin{aligned}
& 2-\sum F_{x}=0 \\
& \therefore P-N_{A}=0 \\
& \therefore P=154.89 \mathrm{~N}
\end{aligned}
$$

## Example: 4

The uniform crate shown in Figure hâs mass of 20 kg . If a force $P=80 \mathrm{~N}$ i applied to the crate, determine if it remains in equilibrium. The coefficient of static friction is $\boldsymbol{\mu}_{s}=\mathbf{0 . 3}$.


## Solution:

From Free-Rody Diagram. the resultant normal force $\boldsymbol{N}_{C}$ must act a distance $x$ from the crate senter line in order to counteract the tipping effect caused by $\boldsymbol{P}$. There are three unknowns, $\boldsymbol{F}, \boldsymbol{N}_{\boldsymbol{C}}$, and $\boldsymbol{x}$, which can be determined strictly from the three equations of equilibrium.


## Equations of Equilibrium.

$$
\begin{aligned}
& \sum F_{x}=0 \\
& \therefore 80 \cos 30^{\circ}-F=0 \\
& \therefore F=80 \cos 30^{\circ}=69.3 \mathrm{~N} \\
& \sum F_{y}=0 \\
& \therefore-80 \sin 30^{\circ}+N_{C}-196.2= \\
& 0 \\
& \sum M_{0}=0 \\
& \therefore 80 \sin 30^{\circ}(0.4)-
\end{aligned}
$$

$80 \cos 30^{\circ}(0.2)+N_{C}(x)=0$
$\therefore N_{C}=236.2 N \quad$ and $\quad \mathrm{x}=-0.00908 \mathrm{~m}=-9.08 \mathrm{~mm}$
Since $x$ is negative it indicates the resultant normal force acts (slightly) to the left of the crate's center line. No tipping will occur since $x<0.4 \mathrm{~m}$. Also, the maximum frictional force which can be developed at the surface of contact is $F_{\max }=\mu_{s} N_{C}=0.3(236.2 \mathrm{~N})=70.9 \mathrm{~N}$.

Since $F=69.3 \mathrm{~N}<70.9 \mathrm{~N}$, the crate will not slip, although it is very close to doing so.

## Example: 5

The uniform 10-kg ladder shown in Figure rests against the smooth wall at $\boldsymbol{B}$, and the end $\boldsymbol{A}$ rests on the rough horizontal plane for which the coefficient of static friction is $\boldsymbol{\mu}_{s}=\mathbf{0 . 3}$. Determine the angle of inclination $\boldsymbol{\theta}$ of the ladder and the normal reaction at $\boldsymbol{B}$ if the ladder is on the verge of slipping.

## Answer:

from the free body diagram, we have
by using the equilibrium conditions
$\sum F_{x}=0 \therefore F_{r}-N_{B}=0$
$\therefore F_{r}=N_{B}$
$\sum F_{Y}=0 \quad \therefore N_{A}-W=0 \quad \therefore N_{A}=W$
From Eqn. (2) $\therefore N_{A}=10 \times 9.81=98.1 \mathrm{~N}$
$\because F_{r}=\mu_{s} N_{A}=0.3 \times 98.1=29.43 \mathrm{~N}$
From Eqn (1) $\therefore N_{B}=F_{r}=29.43 N$
$\sum M_{A}=0 \quad \therefore N_{B} \times h-W \times(L / 2)=0$
$\therefore N_{B} \times h=W \times(L / 2)$

## From the triangle

$h=4 \sin \theta$ and $L=4 \cos \theta$
From Eqn. (3)
$\therefore N_{B} \times 4 \sin \theta=W \times(4 \cos \theta / 2)$
$\therefore 29.43 \sin \theta=49.05 \cos \theta$
$\therefore \theta=\tan ^{-1} \frac{49.05}{29.43}=59.04^{\circ}$

## Example: 6

Determine the distance $s$ to which the $\mathbf{9 0 - k g}$ painter can climb without causing the $4-m$ ladder to slip at its lower end $\boldsymbol{A}$. The top of the $\mathbf{1 5 - k g}$ ladder has a small roller, and at the ground the coefficient of static friction is $\mathbf{0 . 2 5}$. The mass center of the painter is directly above her feet.

## Solution:

$L=\sqrt{4^{2}-1.5^{2}}=3.71 \mathrm{~m}$
$\theta=\cos ^{-1}\left(\frac{1.5}{4}\right)=67.98^{\circ}$


By using the equilibrium conditions, we get
$\sum F_{x}=0 \quad \therefore F-R_{B}=0 \quad \therefore R_{B}=F-\mu R_{A}$

$$
\sum F_{y}=0 \quad \therefore R_{A}-900-150=0 \quad \therefore R_{A}=900+150=1050 \mathrm{~N}
$$

From Eqs. (1) and (2)
$0.25 \times 1050=262.5 N$

$$
\begin{aligned}
& \sum M_{A}=0 \\
& \left(900 \times L_{S}\right)=0
\end{aligned} \quad \because\left(R_{B} \times L\right)-(150 \times 0.75)-
$$

## $(262.5 \times 3.71)-(150 \times 0.75)$

$$
=0.957 \mathrm{~m}
$$

$\because \cos \theta=\frac{L_{S}}{S}$

$\therefore s=\frac{L_{S}}{\cos \theta}=\frac{0.957}{\cos 67.98^{\circ}}=2.55 \mathrm{~m}$

## Problems

1. The $85-\mathrm{lb}$ force $P$ is applied to the $200-\mathrm{lb}$ crate, which is stationary before the force is applied. Determine the magnitude and direction of the friction force $F$ exerted by the horizontal surface on the crate.

2. The $700-\mathrm{N}$ force is applied to the $100-\mathrm{kg}$ block, which is stationary before the force is applied. Determine the magnitude and direction of the friction force $F$ exerted by the horizontal surface on the block.

3. The coefficients of static and kinetic friction between the $100-\mathrm{kg}$ block and the inclined plane are 0.30 and 0.20 , respectively. Determine (a) the friction force $F$ acting on the block when $P$ is applied with a magnitude of 200 N to the block at rest, $(b)$ the force $P$ required to initiate motion up the incline from rest, and $(c)$ the friction force $F$ acting on the block if $P=600 \mathrm{~N}$.

4. The magnitude of force $P$ is slowly increased. Does the homogeneous box of mass $m$ slip or tip first? State the value of $P$ which would cause each occurrence. Neglect any effect of the size of the small feet.


## Chapter (5)

## Area Moments of Inertia

## Solved problems:

## Example:1

Determine the moment of inertia of the cross-sectional area of the channel with respect to the $\boldsymbol{x}_{\boldsymbol{o}}$ and $\boldsymbol{y}_{\boldsymbol{o}}$ axis. Also Determine the moment of inertia of the cross-sectional area of the channel with respect to the $\boldsymbol{x}$ axis.

## Solution:

First, the geometric shape should be divided into a set of geometric shapes.

From the figure can be divided into three parts.

## 1. With respect to the $x_{o}$ axis



## Part: 1

$I_{x_{o_{1}}}=\frac{\mathrm{b}_{1} \mathrm{~h}_{1}^{3}}{12}+\mathrm{A}_{1} d_{1}^{2}=$
$=\frac{30 \times 10^{3}}{12}+(30 \times 10) \times 35^{2}=370000 \mathrm{~mm}^{4}$
Part: 2
$\mathrm{I}_{x_{0_{2}}}=\frac{\mathrm{b}_{2} \mathrm{~h}_{2}^{3}}{12}=\frac{10 \times 60^{3}}{12}=180000 \mathrm{~mm}^{4}$

## Part: 3

$\mathrm{I}_{x_{0}}=\frac{\mathrm{b}_{3} \mathrm{~h}_{3}^{3}}{12}+\mathrm{A}_{3} d_{3}^{2}=$

$=\frac{30 \times 10^{3}}{12}+(30 \times 10) \times 35^{2}=370000 \mathrm{~mm}^{4}$
$I_{x_{o}}=I_{x_{o_{1}}}+\mathrm{I}_{x_{0_{2}}}+\mathrm{I}_{x_{0_{3}}}=370000+180000+370000=920000 \mathrm{~mm}^{4}$
Ans.

## The radius of gyration:

$A_{t}=A_{1}+A_{2}+A_{3}=(30 \times 10)+(60 \times 10)+(30 \times 10)=1200 \mathrm{~mm}^{2}$

$$
k_{x_{o}}=\sqrt{\frac{I_{x_{o}}}{A}}=\sqrt{\frac{920000}{1200}}=27.69 \mathrm{~m}
$$

## 2. With respect to the $y_{o}$ axis

## Part: 1

$I_{y_{o_{1}}}=\frac{\mathrm{h}_{1} \mathrm{~b}_{1}^{3}}{12}=\frac{10 \times 30^{3}}{12}=22500 \mathrm{~mm}^{4}$

## Part: 2

$I_{y_{o_{2}}}=\frac{\mathrm{h}_{2} \mathrm{~b}_{2}^{3}}{12}=\frac{60 \times 10^{3}}{12}=5000 \mathrm{~mm}^{4}$

## Part: 3

$I_{y_{o_{3}}}=\frac{\mathrm{h}_{3} \mathrm{~b}_{3}^{3}}{12}=\frac{10 \times 30^{3}}{12}=22500 \mathrm{~mm}^{4}$
$I_{y_{o}}=I_{y_{o_{1}}}+\mathrm{I}_{y_{0_{2}}}+\mathrm{I}_{y_{0_{3}}}=22500+5000+22500=50000 \mathrm{~mm}^{4}$
Ans.
The radius of gyration:

$$
k_{y_{o}}=\sqrt{\frac{I_{y_{o}}}{A}}=\sqrt{\frac{50000}{1200}}=6.45 \mathrm{~m}
$$

## 3. With respect to the $x$ axis

$I_{x}=I_{x_{o}}+A_{t} d_{x}^{2}$
$\therefore I_{x}=920000+\left(1200 \times 65^{2}\right)=5990000 \mathrm{~mm}^{4}$
Ans.

## The radius of gyration:

$$
k_{x}=\sqrt{\frac{I_{x}}{A}}=\sqrt{\frac{5990000}{1200}}=70.65 \mathrm{~m}
$$

## Example:2

Determine the moment of inertia of the area about the $y$ axis.

## Solution:

A differential element of area that is parallel to the $y$ axis, as shown in Figure, is chosen for integration. Since this element has a thickness $d y$ and intersects the curve at the arbitrary point $(\mathrm{x}, y)$, its area is $d A=\mathrm{y} d x$. Hence, integrating with respect to $y$, from $x=0$ to $x=2 \mathrm{~m}$, yields.
$I_{y}=\int x^{2} d A$

$d A=y d x=\left(4-x^{2}\right) d x$

$$
\begin{aligned}
& I_{y}=2 \int_{0}^{2}\left(4 x^{2}-x^{4}\right) d x=2\left|4 \frac{x^{3}}{3}-\frac{x^{5}}{5}\right|_{0}^{2} \\
&=2\left|4 \frac{2^{3}}{3}-\frac{2^{5}}{5}\right|_{0}^{2}=2 \times \frac{64}{15}=8.53 \mathrm{~m}^{4}
\end{aligned}
$$



## Example:3

Determine the moment of inertia for the rectangular area shown in Fig. 10-5 with respect to (a) the centroidal $x$-axis, (b) the axis $x_{b}$ passing through the base of the rectangle, and (c) the pole or $z$ axis perpendicular to the $\boldsymbol{x}-\boldsymbol{y}$ plane and passing through the centroid $\boldsymbol{O}$.

## Solution:


a. At the centroidal $x$-axis
$d I_{x}=\int y^{2} d A$
Where: $d A=b . d y$
$\therefore I_{x}=\int_{-h / 2}^{h / 2} y^{2} b . d y=b \int_{-h / 2}^{h / 2} y^{2} d y=\left|\frac{y^{3}}{3}\right|_{-h / 2}^{h / 2}=\frac{1}{12} b h^{3}$
$\therefore I_{y}=\frac{1}{12} h b^{3}$


## b. At the $x_{b}$-axis

$\therefore I_{x b}=I_{x}+A d^{2}$
$\therefore I_{x b}=\frac{1}{12} b h^{3}+b . h .\left(\frac{h}{2}\right)^{2}=\frac{1}{3} b h^{3}$
c. At the $z$-axis

$$
\therefore J_{C}=I_{x}+I_{Y}=\frac{1}{12} b h^{3}+\frac{1}{12} h b^{3}=\frac{1}{12}\left[h b^{3}+b h^{3}\right]
$$

## Example:3

Determine the moment of inertia of the cross-sectional area of the channel with respect to the $y$ axis. Determine the radius of gyration of an area about $y$ axis

Solution:
$d I_{y}=x^{2} d A$

$\therefore I_{y}=\int_{x 1}^{x 2} x^{2} d A$

## Part: (1)

$I_{y_{1}}=\frac{1}{12} h b^{3}=\frac{1}{12}(50)(200)^{3}=33333333.3 \mathrm{~mm}^{4}$

Part: (2)
$I_{y_{2}}=\frac{1}{12} h b^{3}=\frac{1}{12}(300)(50)^{3}=3125000 \mathrm{~mm}^{4}$


## Part: (3)

$I_{y_{3}}=\frac{1}{12} h b^{3}=\frac{1}{12}(50)(200)^{3}=33333333.3 \mathrm{~mm}^{4}$

$$
\begin{aligned}
& \therefore I_{y}=I_{y_{1}}+I_{y_{2}}+I_{y_{3}}=33333333.3+3125000+33333333.3=69791666.6 \mathrm{~mm}^{4} \\
&=69.8 \times 10^{6} \mathrm{~mm}^{4}
\end{aligned}
$$

The radius of gyration:

$$
k_{y}=\sqrt{\frac{I_{y}}{A}}=\sqrt{\frac{69791666.6}{35000}}=44.65 \mathrm{~m}
$$

## Example:4

Determine the moments of inertia of the rectangular area about the centroidal $\boldsymbol{x}_{0^{-}}$and $\boldsymbol{y}_{0}$-axes, the centroidal polar axis $\mathbf{z}_{0}$ through $\boldsymbol{C}$, the $x$-axis, and the polar axis $z$ through $\boldsymbol{O}$.

## Solution:

For the calculation of the moment of inertia $\boldsymbol{I}_{x 0}$ about the $\boldsymbol{x}_{0}$ axis, a horizontal strip of area $\boldsymbol{b} \boldsymbol{d y}$ is chosen so that all elements of the strip have the same $y$-coordinate. Thus,
$I_{x_{0}}=\int y^{2} d A$


Where a horizontal strip of area $\boldsymbol{d} \boldsymbol{A}=\boldsymbol{b} \boldsymbol{d} \boldsymbol{y}$
$I_{x_{0}}=\int_{-h / 2}^{+h / 2} y^{2} b d y=b\left|\frac{y^{3}}{3}\right|_{-h / 2}^{h / 2}=\frac{1}{12} b h^{3}$
By interchange of symbols, the moment of inertia about the centroidal $y_{0}$-axis is
$I_{y_{0}}=\int_{-b / 2}^{+b / 2} x^{2} h d x=h\left|\frac{x^{3}}{3}\right|_{-b / 2}^{b / 2}=\frac{1}{12} h b^{3}$
The centroidal polar moment of inertia is
$I_{z_{0}}=I_{x_{0}}+I_{y_{0}}$
$=\frac{1}{12} b h^{3}+\frac{1}{12} h b^{3}=\frac{1}{12} b h\left(b^{2}+h^{2}\right)=\frac{1}{12} A\left(b^{2}+h^{2}\right)$
By the parallel-axis theorem the moment of inertia about the $\boldsymbol{x}$-axis is
$I_{x}=I_{x_{0}}+A d^{2}$
$=\frac{1}{12} b h^{3}+(b \cdot h)\left(\frac{h}{2}\right)^{2}=\frac{1}{3} b h^{3}=\frac{1}{3} A h^{2}$
We also obtain the polar moment of inertia about $O$ by the parallel-axis theorem, which gives us
$I_{z}=I_{Z_{0}}+A d^{2}$
$=\frac{1}{12} A\left(b^{2}+h^{2}\right)+(b . h)\left[\left(\frac{h}{2}\right)^{2}+\left(\frac{h}{2}\right)^{2}\right]=\frac{1}{3} b h\left(b^{2}+h^{2}\right)=\frac{1}{3} A\left(b^{2}+h^{2}\right)$

